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Regular black holes without mass-inflation instability and gravastars from modified gravity

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Singularities and regular black holes

- Vacuum black holes in GR have singularities

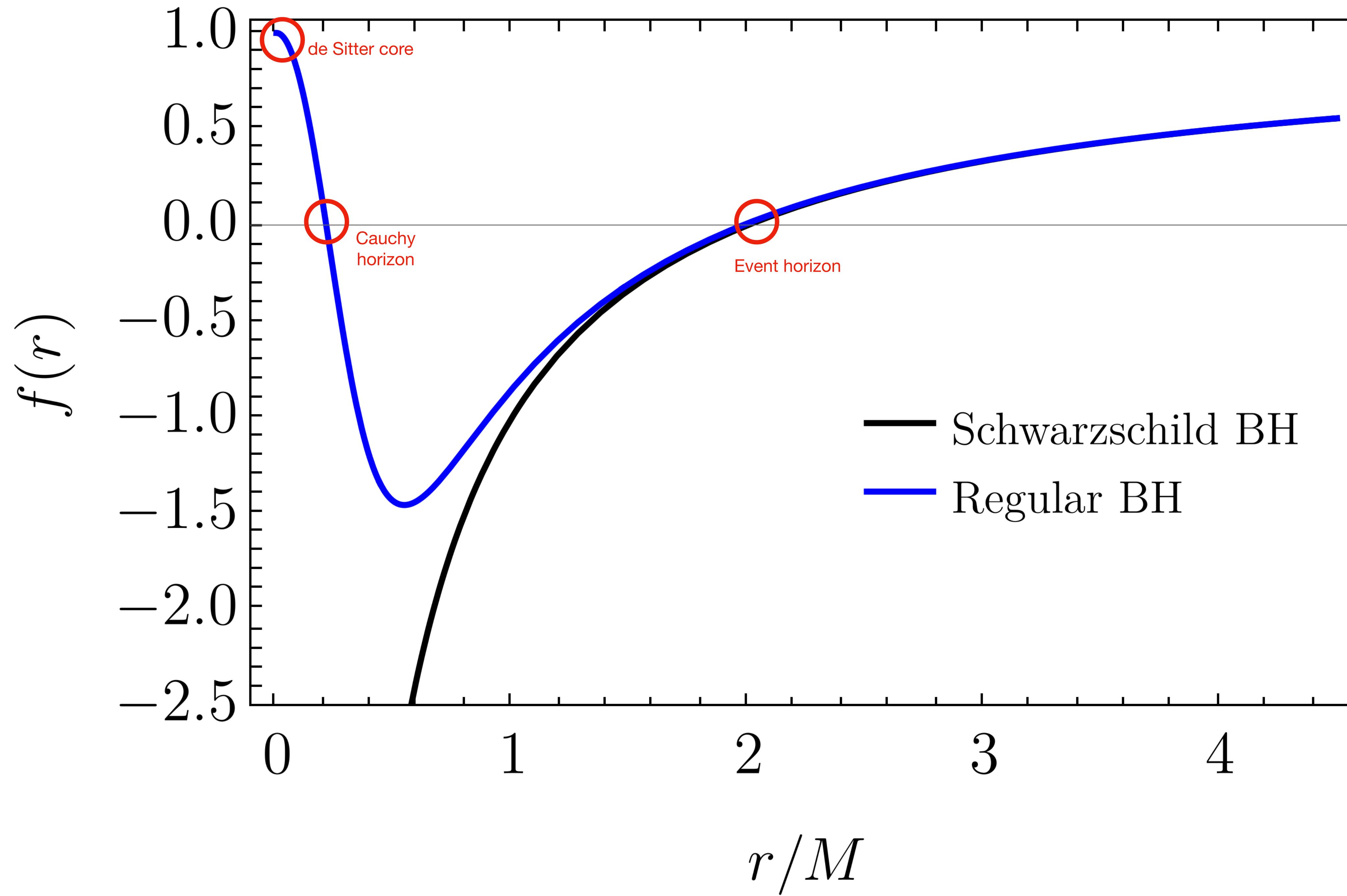
$$ds^2 = -f(r)dt^2 + f(r)^{-1}dr^2 + r^2d\Omega^2, \quad f(r) = 1 - \frac{2M}{r}$$

$$R_{\mu\alpha\nu\beta}R^{\mu\alpha\nu\beta} = \frac{48M^2}{r^6}, \quad \text{diverges as } r \rightarrow 0$$

- Regular black holes are a proposed solution to this problem
- A metric description remains valid in the core, and is non-singular everywhere
- Example: Hayward metric [Hayward, 2006]

$$f(r) = 1 - \frac{2Mr^2}{r^3 + 2M\ell^2}, \quad f(r) \sim 1 + \mathcal{O}(r^2) \quad \text{near } r = 0$$

Length scale associated with new physics



There are challenges with regular black holes

The main ones

1. Regular BHs of generic mass are *rarely* found as solutions to the field equations of a theory **[Challenge 1]**
2. The Cauchy horizon is **unstable**: if its surface gravity is non-zero, perturbations grow exponentially in a phenomenon known as **mass-inflation** **[Challenge 2]**

Here we discuss a framework where both challenges are addressed

The theory

Motivation and brief derivation

- In four-dimensions, under reasonable assumptions, GR is unique (Lovelock's theorem)
- In higher-dimensions the first non-trivial correction to GR comes in the form of a Gauss-Bonnet term, $\mathcal{G} = R_{\mu\alpha\nu\beta}R^{\mu\alpha\nu\beta} - 4R_{\mu\nu}R^{\mu\nu} + R^2$
- In 4D the GB term is topological/trivial
- Much work has been done in recent years to “dimensionally regularize” the GB term to get a 4D-GB theory of gravity [see the review PF et. al, 2022]

$$\sim \lim_{D \rightarrow 4} \int d^D x \sqrt{-g} \frac{\mathcal{G} - \text{counterterms}}{D - 4}$$

The theory

Motivation and brief derivation

- In performing one of such procedures, by considering different connections with Weyl vectors, we arrive at the following 4D-GB theory

$$S = \int d^4x \sqrt{-g} \left[R + \ell^2 (\mathcal{L}[A] - \mathcal{L}[B]) \right]$$

Length scale associated with new-physics

where $\mathcal{L}[X] = 4G^{\mu\nu}X_\mu X_\nu + 8X_\mu X^\mu \nabla_\nu X^\nu + 6(X_\mu X^\mu)^2$

- It is a 4D theory, with two Proca fields, and second-order equations of motion
- For details see [Charmousis et. al 2025, Eichhorn and PF 2025]

Black hole solutions and regularity

$$A_\mu dx^\mu = a(r) \left(dt^2 + \frac{dr^2}{f(r)} \right) \quad B_\mu dx^\mu = b(r) \left(dt^2 + \frac{dr^2}{f(r)} \right)$$

$$a(r) = \frac{r - rf(r) - \cancel{q_a/2}}{2r^2}, \quad b(r) = \frac{r - rf(r) - \cancel{q_b/2}}{2r^2}$$

$$f(r) = 1 - \frac{2Mr^2 + \ell^2(q_a^2 - q_b^2)/(4r)}{r^3 + (q_a - qb)\ell^2}$$

Regularity is achieved if $q_b = -q_a$

$$f(r) = 1 - \frac{2Mr^2}{r^3 + 2q_a\ell^2}$$

Usual Hayward metric is recovered
when $q_a = M$

Challenge 1

We have a theory where black holes of all masses can be regular

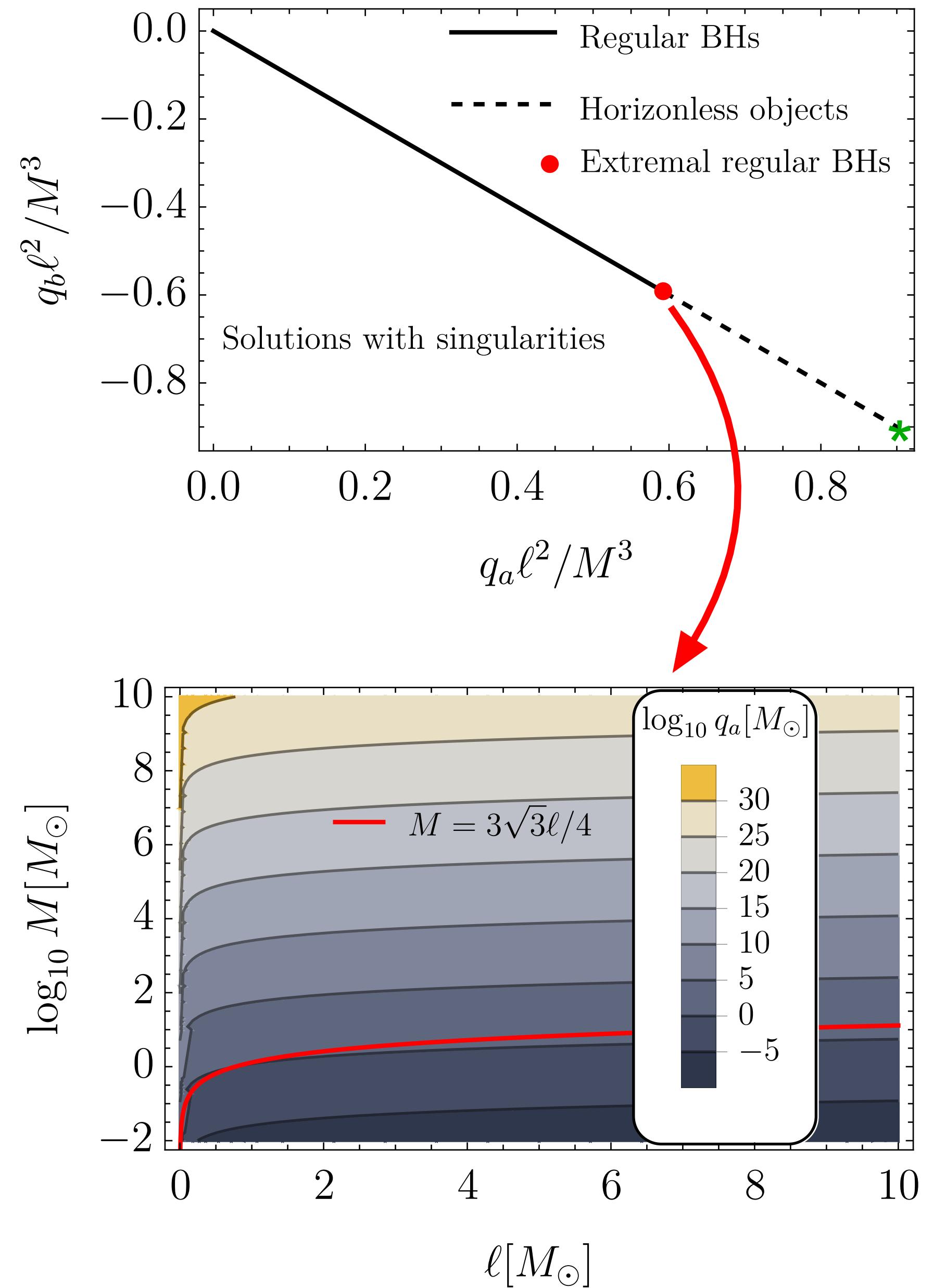
What about Challenge 2? (Stability against mass-inflation instability)

Mass-inflation instability

- A black hole with a non-extremal Cauchy horizon generally suffers from mass-inflation instability
- No instability only if the Cauchy horizon is extremal (zero surface gravity)
[Carballo-Rubio et. al, 2022]
- Regular extremal black holes of all masses are possible in our theory if

$$q_a = \frac{16M^3}{27\ell^2}$$

$$f(r) = 1 - \frac{2Mr^2}{r^3 + \frac{32}{27}M^3}$$



Challenge 2

We have a theory where black holes of all masses can be regular and free from mass-inflation instability

Bonus: Gravastars

- Gravastars are proposed alternatives to regular black holes [Mazur and Mottola 2001]
- Regular BHs smoothly interpolate between a de Sitter core and a Schwarzschild exterior
- Gravastars are defined by a sharp transition

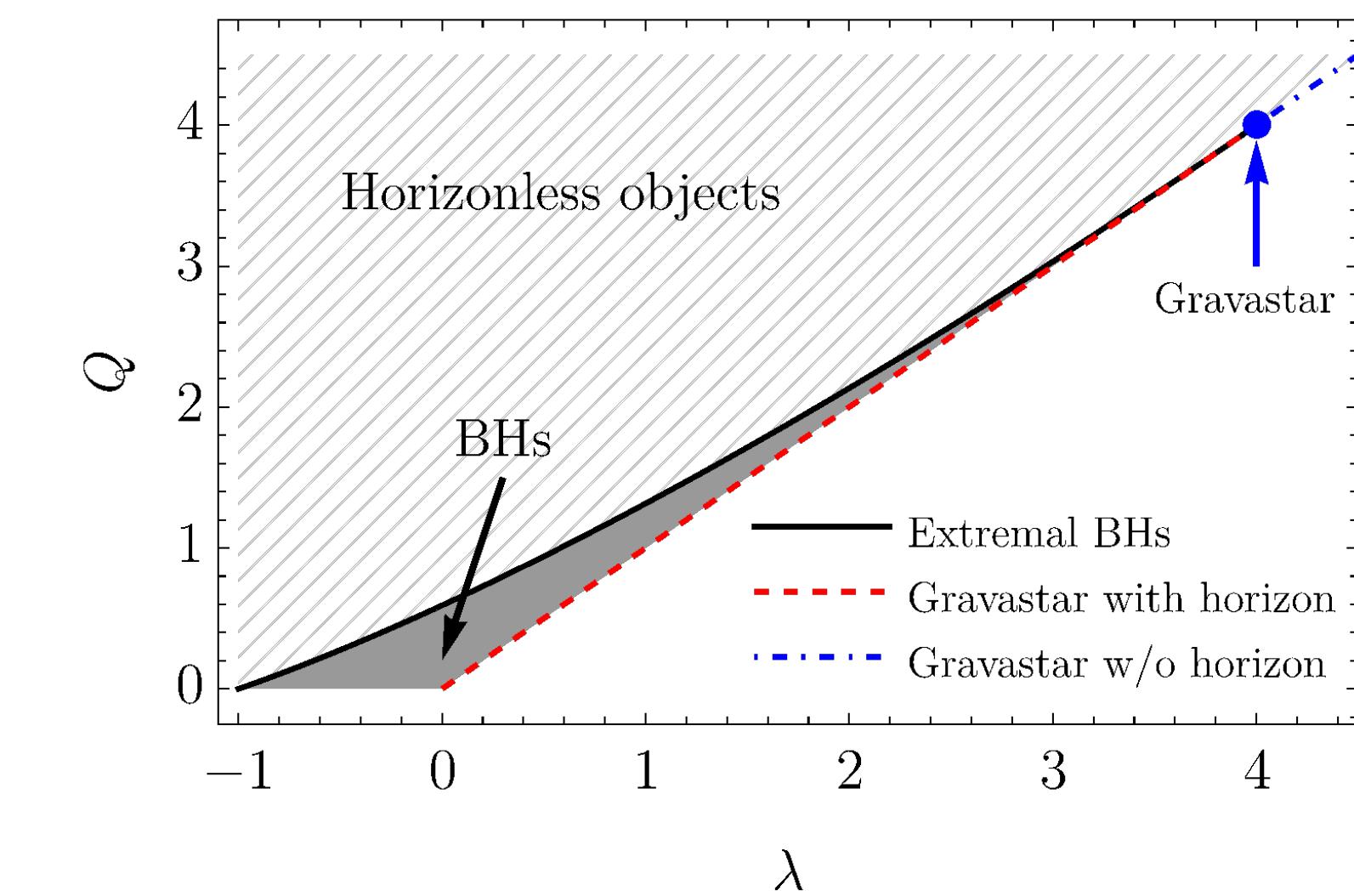
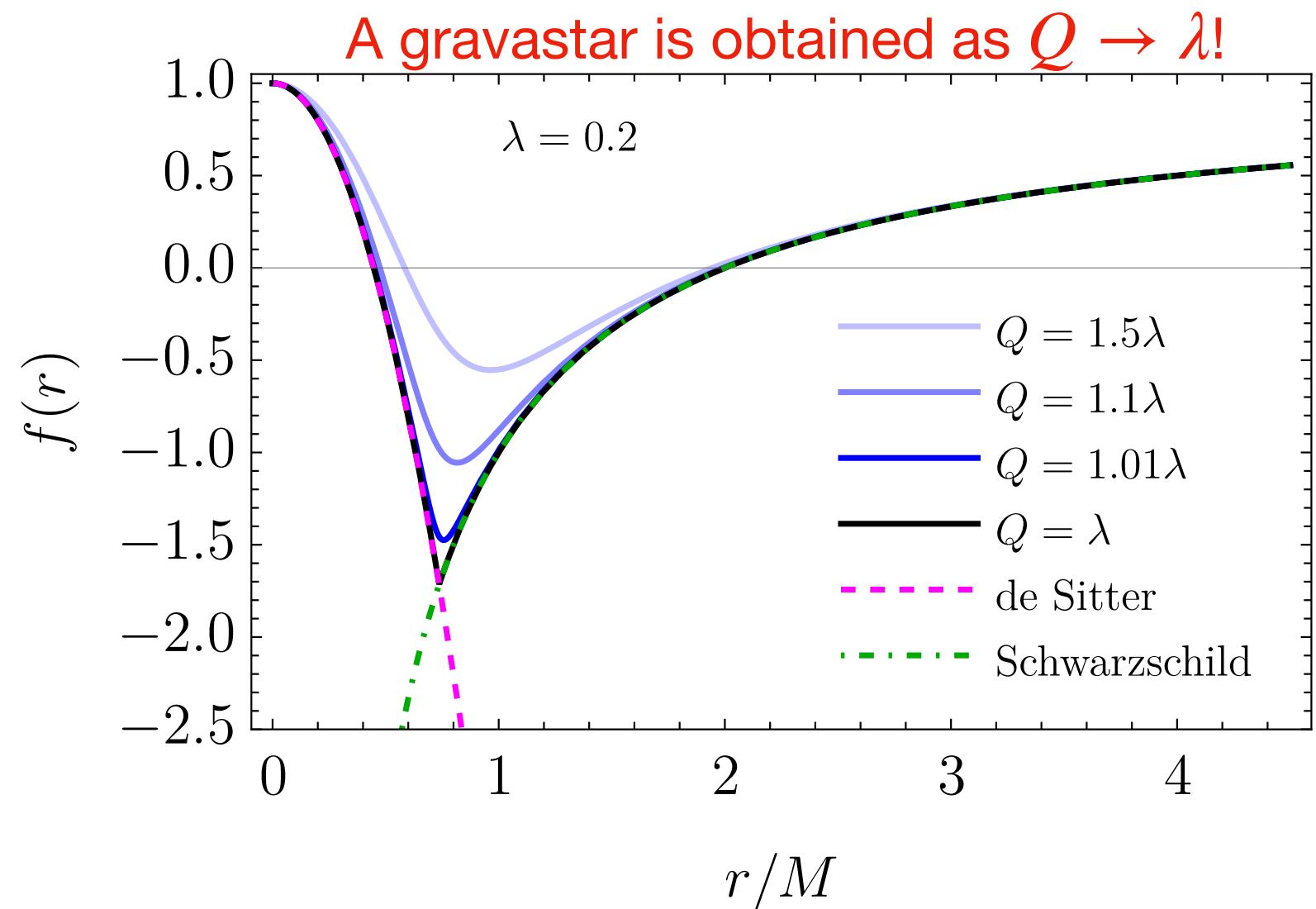
$$f(r) = \begin{cases} 1 - \frac{\Lambda_{\text{eff}}}{3} r^2, & r \leq r_{\text{trans}} \\ 1 - \frac{2M}{r}, & r > r_{\text{trans}} \end{cases}$$

- Thin shell is needed
- No known theories where gravastars arise as solutions

Bonus: Gravastars

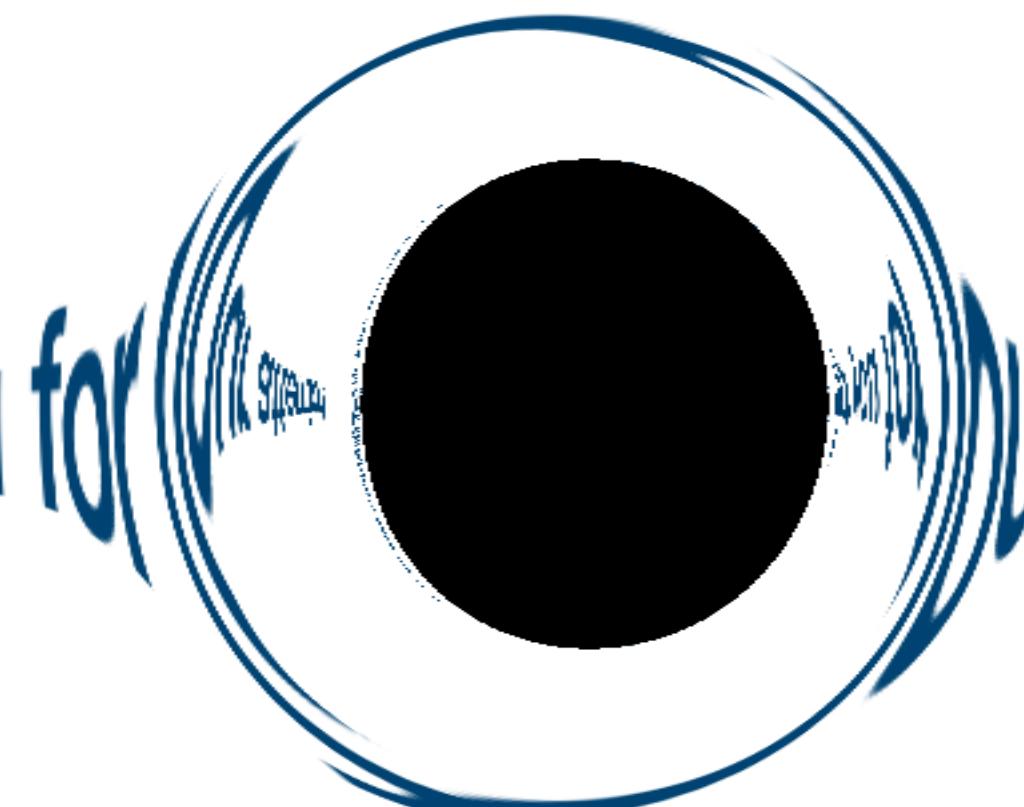
$$S = \int d^4x \sqrt{-g} \left[R + \ell^2 (\mathcal{L}[A] - \kappa \mathcal{L}[B]) \right]$$

$$f(r) = 1 - \frac{MQ}{r\lambda} - \frac{r^2}{2\lambda M^2} \left(1 - \sqrt{\left(1 + \frac{2M^3Q}{r^3} \right)^2 - \frac{8\lambda M^3}{r^3}} \right), \quad \text{where} \quad Q = \frac{q_a \ell^2 (1 + \sqrt{\kappa})}{2M^3}, \quad \lambda = \frac{(1 - \kappa) \ell^2}{M^2}$$



Open questions

- Are the theories healthy? No canonical kinetic term for the Proca fields. Is this an issue?
- Can the values of q_a that eliminate mass-inflation be achieved dynamically? Numerical simulations needed
- Are these regular BHs good dark matter candidates? If BHs of all masses are extremal, do we solve the information paradox?
- Are the horizonless compact objects (obtained for large q_a) stable against the light-ring instability?
- Sizeable deviations from GR at all masses. Observational signatures?
- ...



Thank you for your attention!