

Thermodynamics of hot curved space in the canonical ensemble and black hole nucleation

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Cavity surrounded by a heat reservoir

Euclidean path integral approach

Gibbons and Hawking¹: $Z = \int Dg e^{-I[g]}$, $Z_{\text{0loop}} = e^{-I[g_{\text{cl}}]} = e^{-\beta F}$;

Gross, Perry and Yaffe²: Black hole nucleation from hot flat space;

York's realization³: Finite cavity.

Entropy of self-gravitating radiation fluid

Sorkin, Wald and Zhang⁴: radiation as perfect fluid inside a spherical cavity.

¹G. W. Gibbons and S. W. Hawking, Phys. Rev. D **15**, 2752 (1977).

²D. J. Gross, M. J. Perry and L. G. Yaffe, Phys. Rev. D **25**, 330 (1982).

³J. W. York, Phys. Rev. D **33**, 2092 (1986).

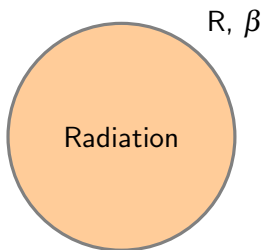
⁴R. D. Sorkin, R. M. Wald and Z. J. Zhang, Gen. Rel. Grav. **13**, 1127 (1981).

Self-gravitating fluid inside a cavity

The setup

The action

$$I = -\frac{1}{16\pi l_{\text{p}}^2} \int \sqrt{g} R d^4x + \int \mathcal{F}_{\text{m}} d^4x - \frac{1}{8\pi l_{\text{p}}^2} \int (K - K_0) \sqrt{\gamma} d^3x$$



Spherical symmetry

$$ds^2 = b(y)^2 d\tau^2 + a(y)^2 dy^2 + r^2 d\Omega^2$$

Boundary conditions

Regularity: $b(0)$ finite, $\frac{b'}{a}|_{y=0} = 0$, $r(0) = 0$;

Heat reservoir: $\beta = 2\pi b(1)$, $r(1) = R$;

Perfect fluid: $\varepsilon = a T_{\text{loc}}^4$;

Local temperature: $T_{\text{loc}} = \frac{1}{2\pi b(y)}$
(ensures the conservation of the stress-tensor).

Self-gravitating fluid inside a cavity

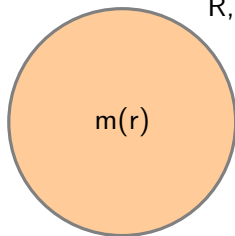
Reduced action, equilibrium and stability

Imposing the Hamiltonian constraints, the action becomes ($f = 1 - \frac{2m(r)}{r}$)

$$I(R, \beta; m(r)) = \frac{R\beta}{l_p^2} \left(1 - \sqrt{f}\right) - 4\pi \int_0^R s r^2 f^{-1/2} dr ; m(r) = 4\pi \int_0^r \varepsilon r^2 dr$$

Equilibrium equation ($\delta I = 0$)

$$R, \beta \quad \frac{1}{\varepsilon + p} \frac{dp}{dr} + \frac{l_p^2}{r^2 f} (4\pi r^3 p + m) = 0 \quad (\text{TOV equation})$$



Stability conditions ($\delta^2 I > 0$)

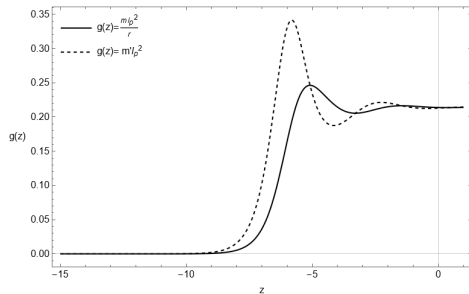
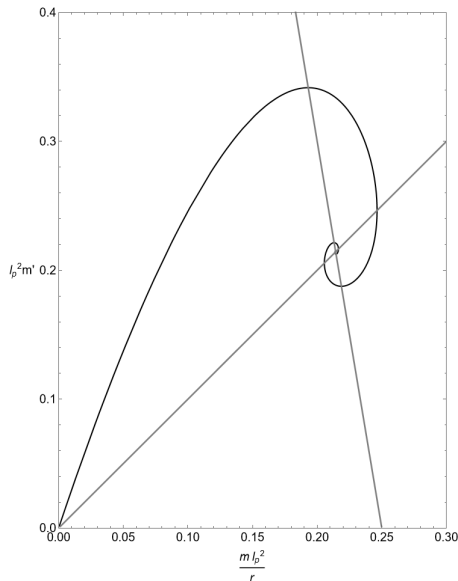
$$\hat{L}\delta m = -\chi \frac{\delta m \pi T_{\text{loc}}}{r^2 f^{\frac{3}{2}} (\varepsilon + p)} ; \delta m'(R) = 0$$

Stable if $\chi > 0$ for all modes δm

$$\hat{L}\delta m = \left(\frac{\delta m'}{4\pi r^2 \sqrt{f} T_{\text{loc}}^2} \frac{dT_{\text{loc}}}{d\varepsilon} \right)' + \frac{2l_p^2}{r T_{\text{loc}}^2 f^{\frac{3}{2}}} \left(\frac{T_{\text{loc}}}{2r} - \frac{dT_{\text{loc}}}{dr} \right)$$

Self-gravitating fluid inside a cavity

Equilibrium solutions



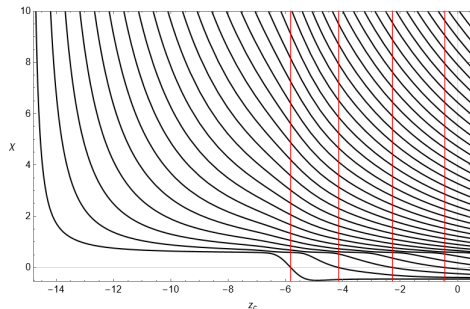
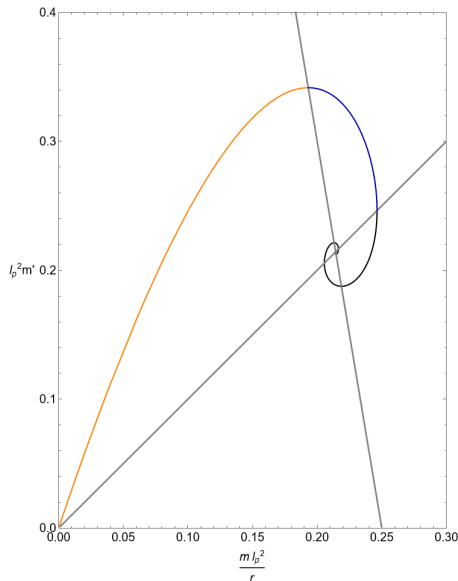
$$r \frac{d(m')}{dr} = \frac{2m'(1 - 4\frac{l_p^2 m}{r} - \frac{2}{3}l_p^2 m')}{(1 - 2\frac{l_p^2 m}{r})}$$

$$r \frac{d(m/r)}{dr} = m' - \left(\frac{m}{r}\right)$$

$$z \propto \log(r)$$

Self-gravitating fluid inside a cavity

Thermodynamic and mechanical stability



Thermodynamic and mechanical stability (orange)

$$\hat{L}\delta m = -\chi \frac{\delta m \pi T_{\text{loc}}}{r^2 f^{\frac{3}{2}}(\varepsilon + p)} ; \delta m'(R) = 0$$

Mechanical stability (blue)

$$\delta m(R) = 0$$

Self-gravitating fluid inside a cavity

Thermodynamics

Entropy

$$S = \frac{2(4\pi a)^{\frac{1}{4}}}{9} \left(\frac{R}{l_p} \right)^{\frac{3}{2}} \frac{l_p^2 (m' + \frac{m}{R})}{(l_p^2 m')^{\frac{1}{4}} \sqrt{1 - \frac{2l_p^2 m}{R}}}$$

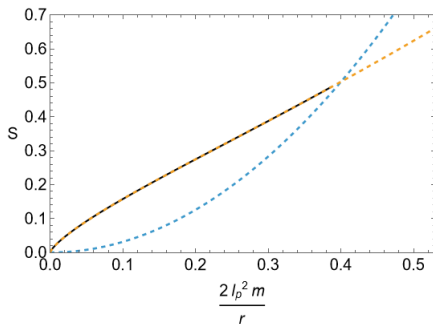
$$R = l_p:$$

$$S \propto \left(\frac{2l_p^2 m}{R} \right)^{0.755} + 0.5 \left(\frac{2l_p^2 m}{R} \right)^{2.24}$$

$$S_{\text{bh}} = \pi \left(\frac{r_+}{R} \right)^2$$

Energy

$$E = \frac{R}{l_p^2} \left(1 - \sqrt{1 - \frac{2l_p^2 m}{R}} \right)$$

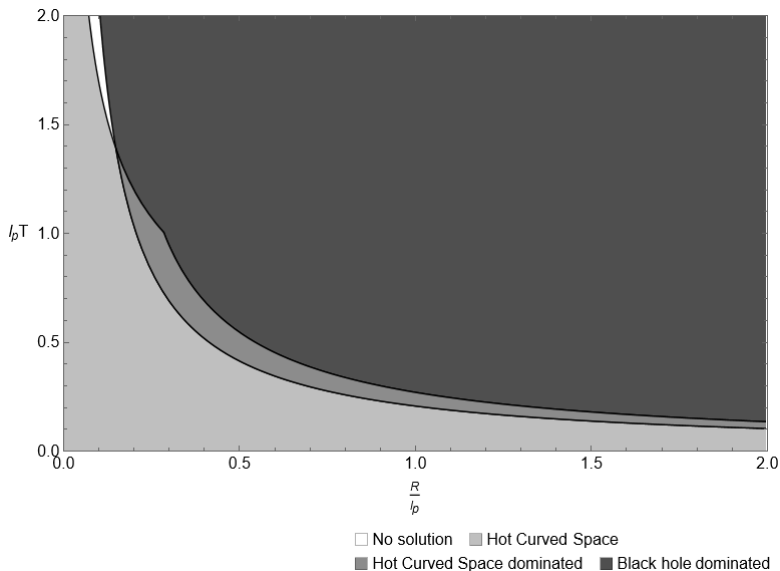


Black - radiation entropy;

Orange - fit to radiation entropy;

Blue - black hole entropy;

Phase diagram Self-gravitating radiation vs Black hole



Construction of the canonical ensemble of a radiation fluid inside a cavity:

- There is a stable solution that exists for a maximum compactness ($\frac{l_p^2 m}{R} \approx 0.19$) and temperature;
- Thermodynamic and mechanic stability is represented by one single condition;
- Entropy goes as $\left(\frac{l_p^2 m}{R}\right)^{0.755}$ at first order;
- There is black hole nucleation;
- Picture becomes intricate beyond the Planck regime;