

A greybody-factor approach to modeling black hole ringdown



XVIII
BLACK HOLES
WORKSHOP

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Lisbon 19/12/2025

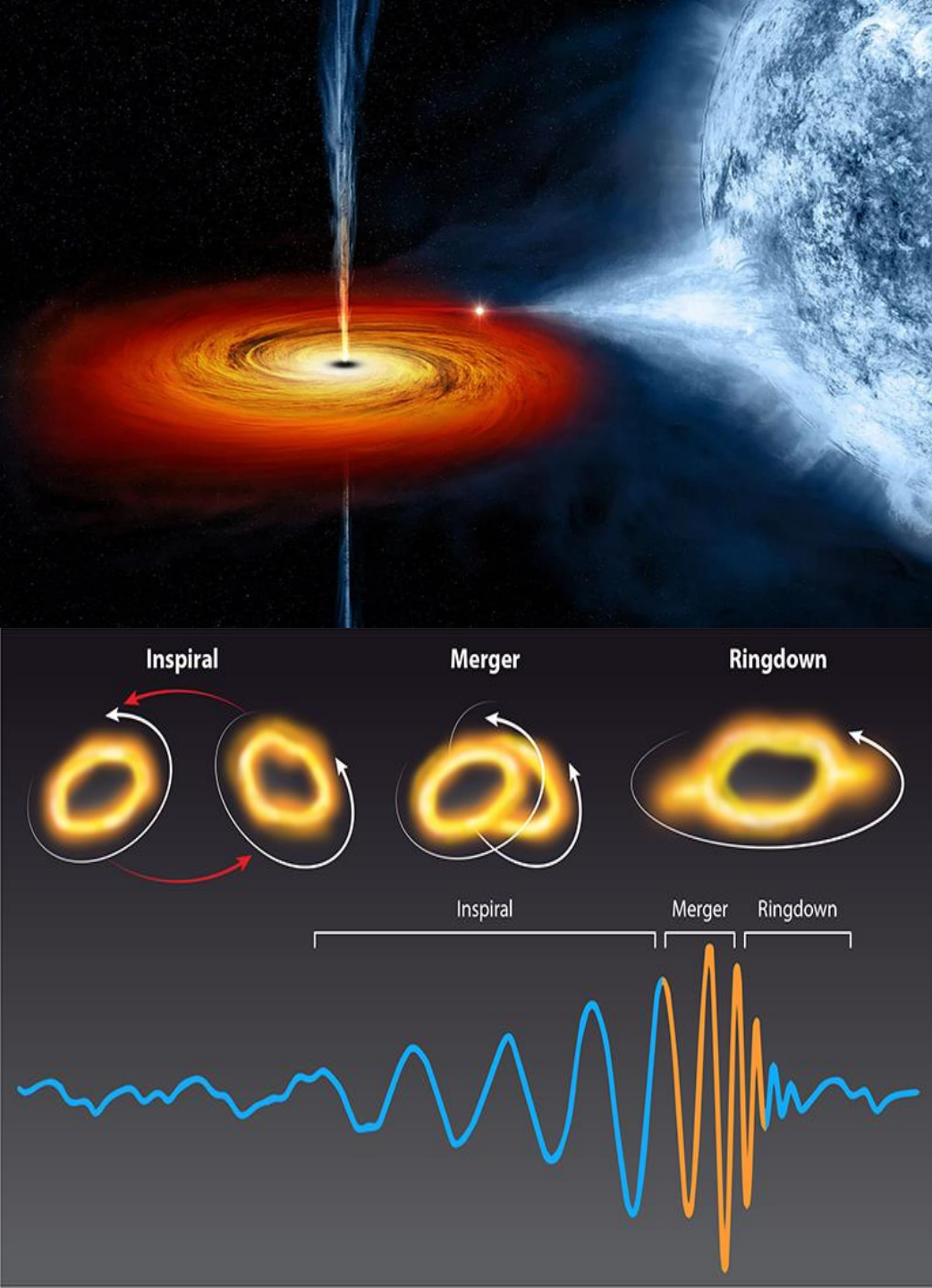
<https://arxiv.org/pdf/2512.15877>

In collaboration with Sophia Yi,
Emanuele Berti and Paolo Pani

OVERVIEW

1. **Ringdown signals:** greybody factors vs quasinormal modes.

2. **Comparable mass mergers:** fitting frequency domain ringdown amplitude with greybody factors.

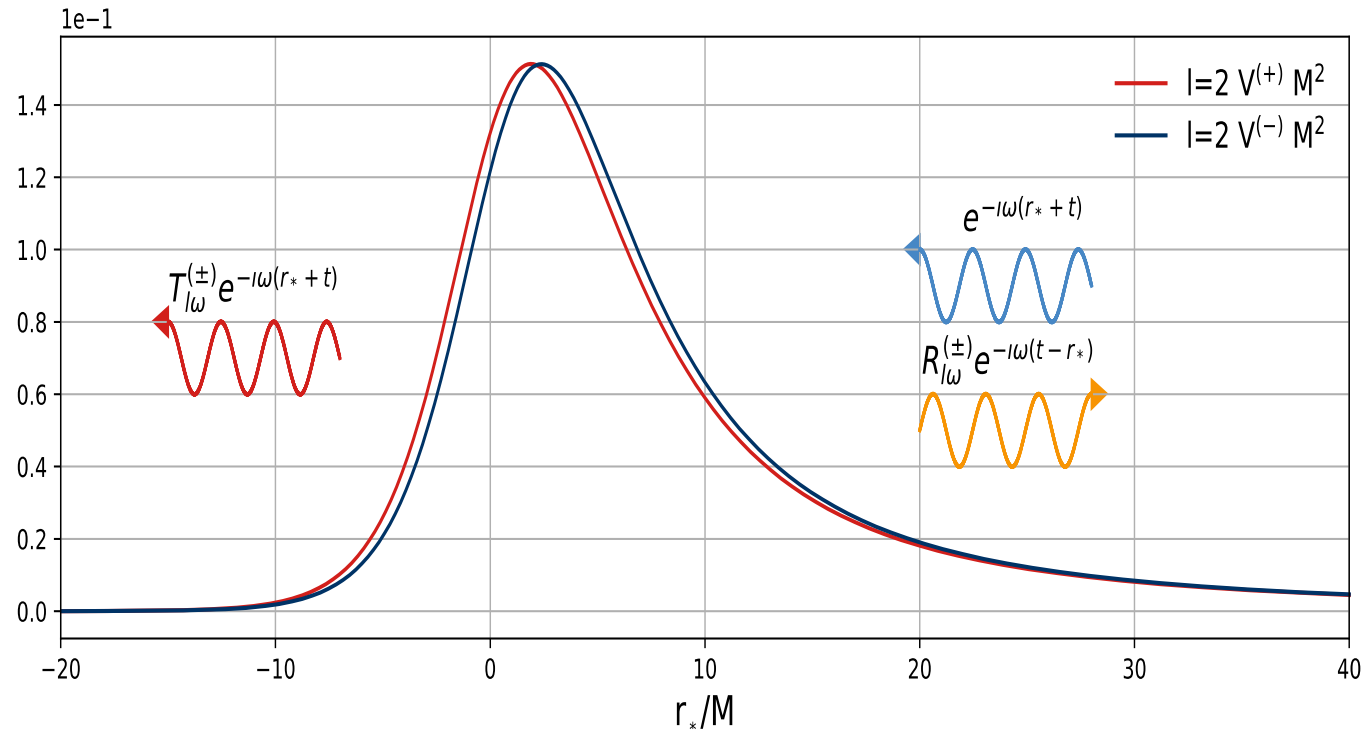


QUASINORMAL MODES vs GREYBODY FACTORS

Perturbations of a Schwarzschild black hole can be described through **Regge-Wheeler** and **Zerilli** equations.

$$\left(\frac{d^2}{dr_*^2} - V_l^{(\pm)}(r) + \omega^2 \right) Z_l^{(\pm)}(r_*) = 0$$

Greybody factors (GFs) describe how the black hole reacts to a scattering process.



$$Z_l^{(\pm)}(r_*) \begin{cases} A_{l\omega}^{(\pm),out} e^{i\omega r_*} + A_{l\omega}^{(\pm),in} e^{-i\omega r_*} & r_* \rightarrow \infty \\ e^{-i\omega r_*} & r_* \rightarrow -\infty \end{cases}$$

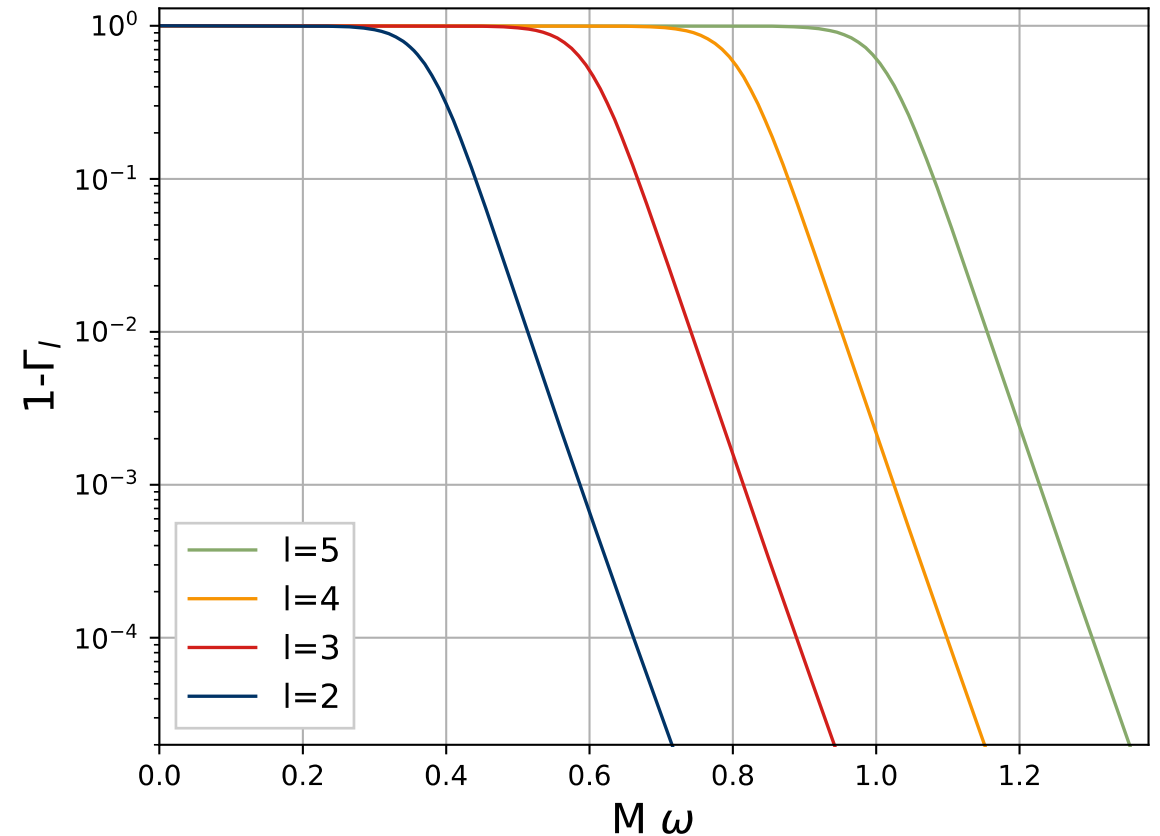
GREYBODY FACTORS

From this, one can define black hole reflectivity and black hole transmission amplitude (i.e. the **greybody factor** $\Gamma_l(\omega)$).

GFS AND REFLECTIVITY

$$\Gamma_l(\omega) = \left| T_{l\omega}^{(\pm)} \right|^2 = \left| \frac{1}{A_{l\omega}^{(\pm),in}} \right|^2$$

$$\mathcal{R}_l(\omega) = \left| R_{l\omega}^{(\pm)} \right|^2 = \left| \frac{A_{l\omega}^{(\pm),out}}{A_{l\omega}^{(\pm),in}} \right|^2 = 1 - \Gamma_l(\omega)$$



SOME RESULTS

- GFs have been showed to be imprinted the frequency domain ringdown amplitude in the point particle infall scenario for BHs.
[see N.Oshita <https://arxiv.org/pdf/2309.05725>, **RFR**, K. Destounis and P. Pani <https://arxiv.org/abs/2406.01692>]
- GFs have been showed to be stable under small perturbations of the wave equation (differently from QNMs that are highly unstable)
[see **RFR**, K. Destounis and P. Pani <https://arxiv.org/abs/2406.01692>, N. Oshita + <https://arxiv.org/pdf/2406.04525>]
- The GFs framework has been showed to work also for exotic compact objects and wormholes.
[see **RFR**, S. Biswas, S. Chackraborty and P. Pani <https://arxiv.org/pdf/2501.16433>]

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HOW TO MODEL GRAVITATIONAL WAVES WITH GFS

We all know that the time domain ringdown can be described as a superposition of quasinormal oscillations.

In the greybody factor framework however we focus on **frequency domain**.

$$H_{\ell m}(\omega) = |h_{\ell m}(\omega)| \quad \text{with} \quad h_{\ell m}(\omega) = \frac{1}{2\pi} \int dt \, h_{\ell m}(t) e^{-i\omega t}$$

The idea is really simple. The reflectivity is imprinted in the Fourier transform in the sense:

$$H_{\ell m}(\omega) \propto \mathcal{R}_{\ell m}(M\omega, \chi)$$

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Fourier Transform
of the signal

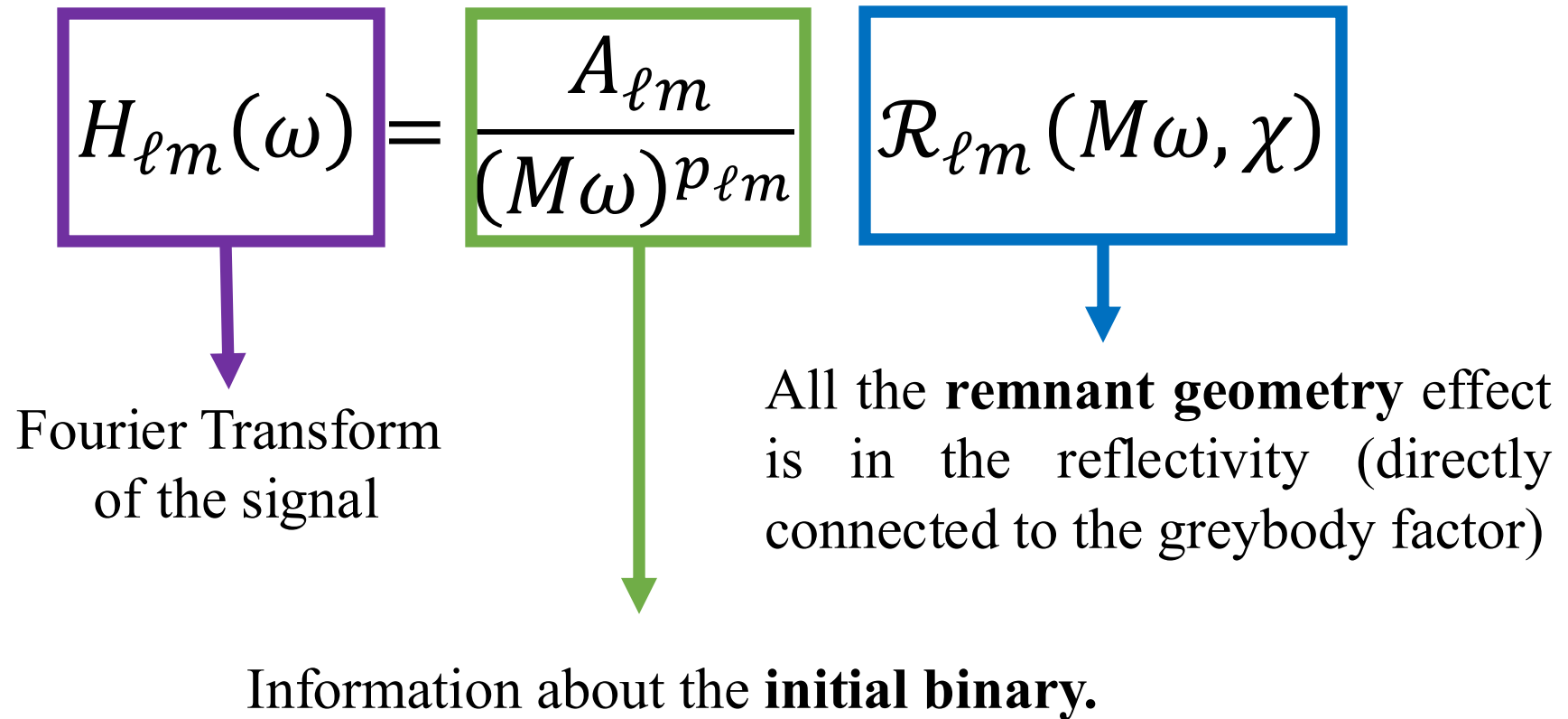
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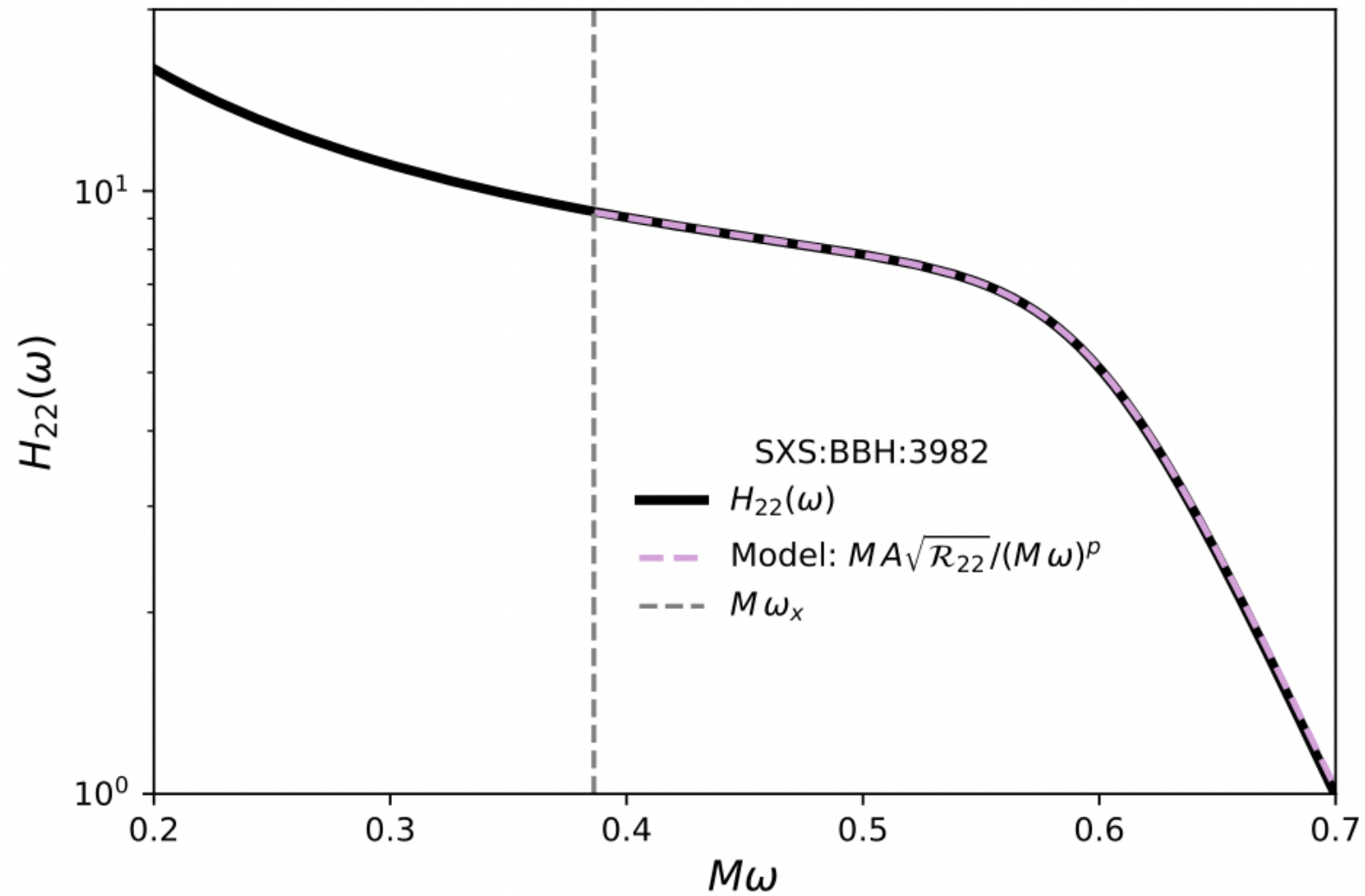
Fourier Transform
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All the **remnant geometry** effect
is in the reflectivity (directly
connected to the greybody factor)

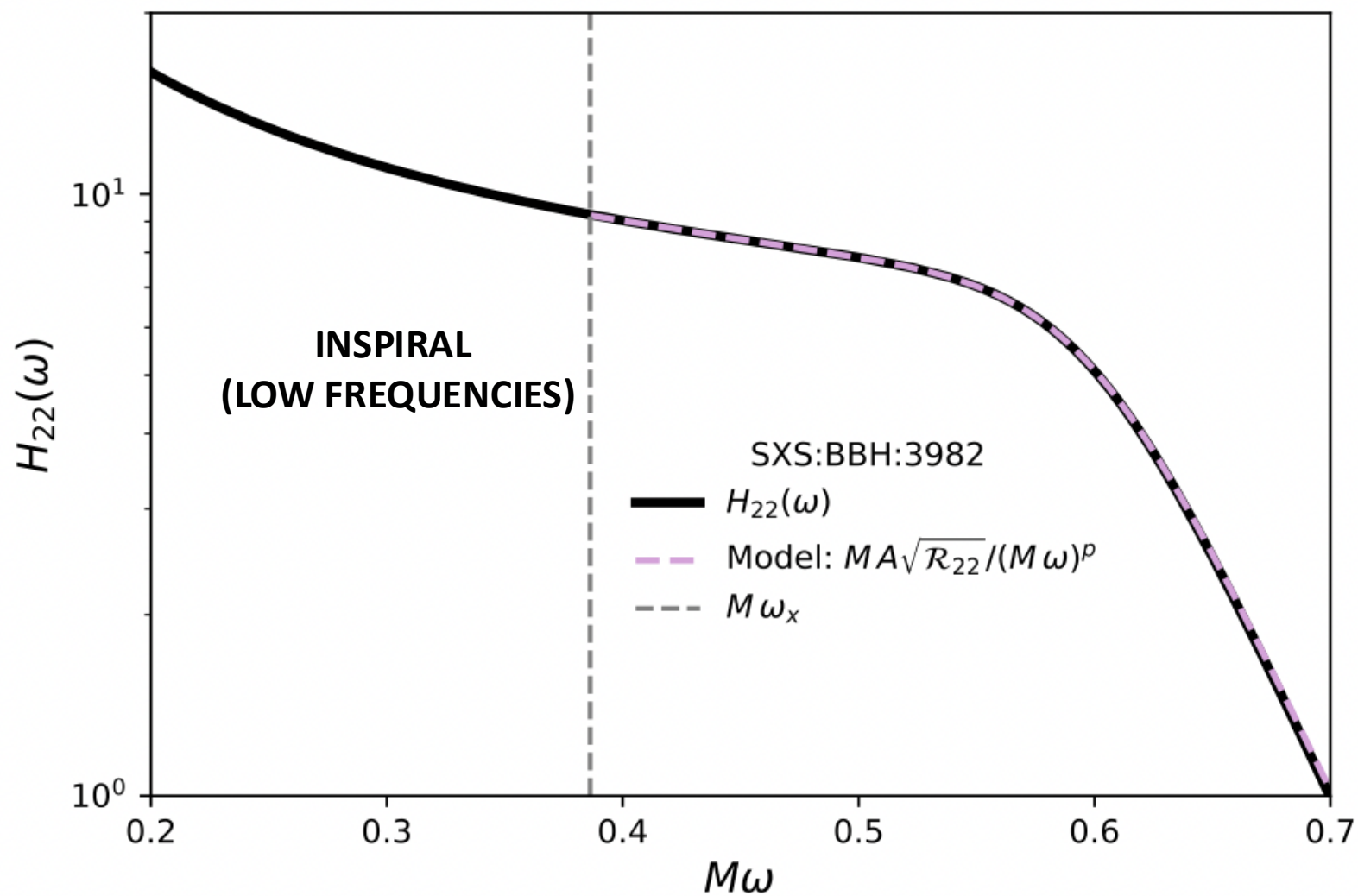
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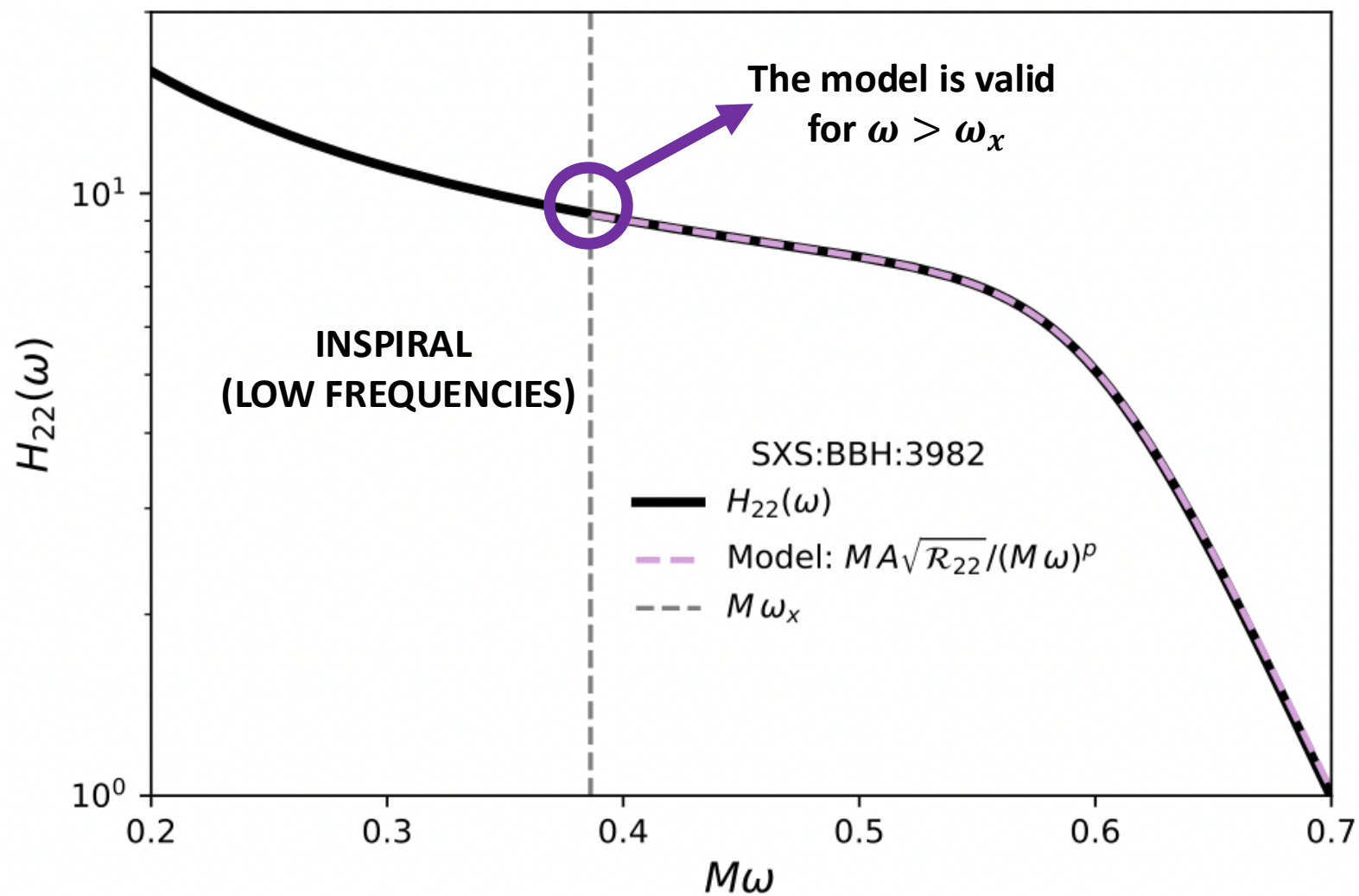
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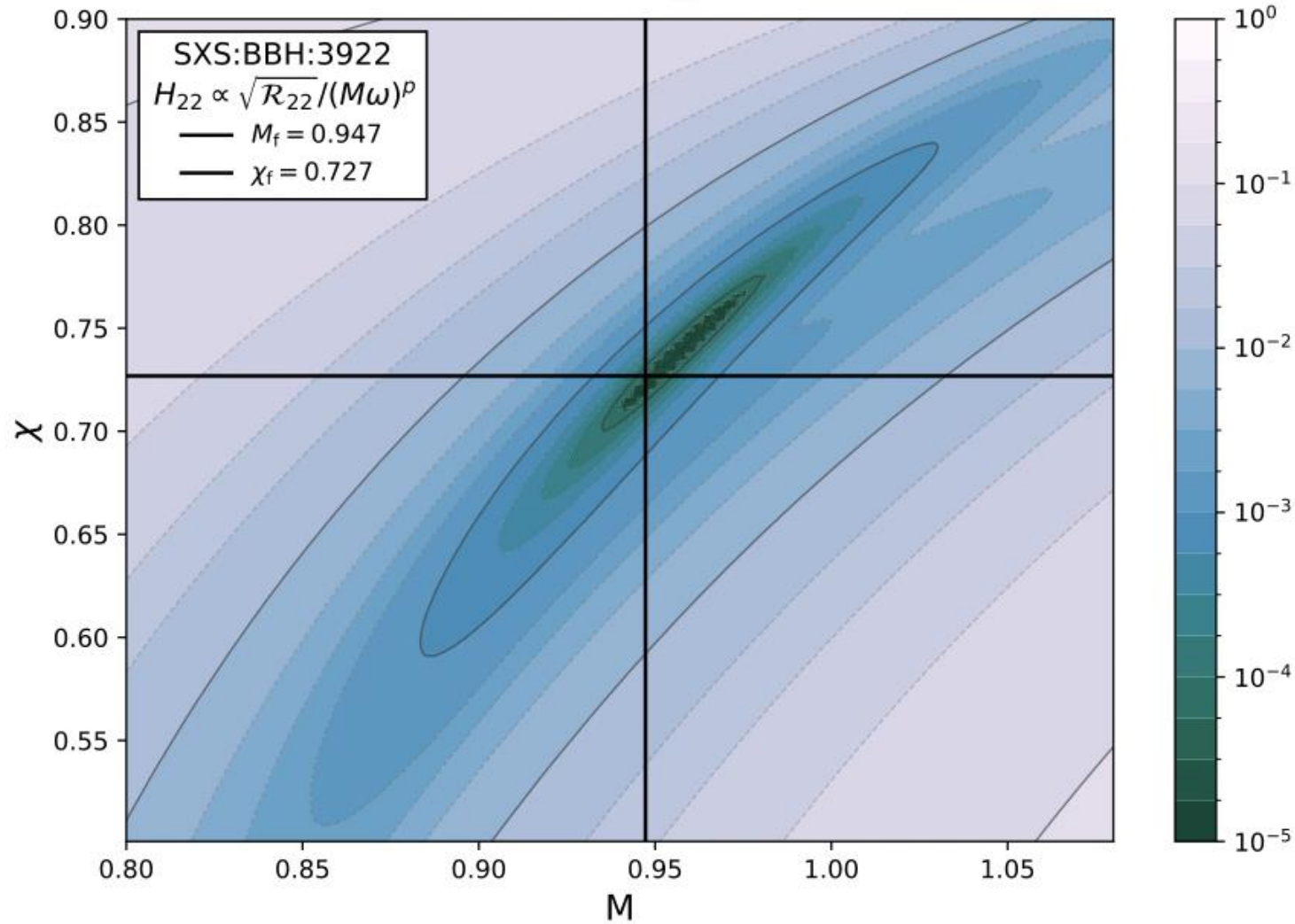
HOW TO MODEL GRAVITATIONAL WAVES WITH GFS



HOW TO MODEL GRAVITATIONAL WAVES WITH GFS



I. DETERMINATION OF THE REMNANT PARAMETERS



$$\mathcal{M} = 1 - \frac{\langle H^{\text{data}} | H^{\text{model}} \rangle}{\sqrt{\langle H^{\text{data}} | H^{\text{data}} \rangle \langle H^{\text{model}} | H^{\text{model}} \rangle}}$$

The true values of the remnant parameters lie in the minimal mismatch region, achieving mismatches of $\mathcal{O}(10^{-5})$.

II. FITTING $A_{\ell m}$ and $p_{\ell m}$ WITH BINARY PARAMETERS

Information about the **initial binary** is contained in $A_{\ell m}$ and $p_{\ell m}$.

$$A_{\ell m} = A_{\ell m}(\mathbf{m}_1, \mathbf{m}_2, \vec{\chi}_1, \vec{\chi}_2)$$

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We fitted a subset of the SXS catalog simulations as polynomials of the binary parameters, using the following parametrization

$$\chi_+ = \frac{q\chi_1 + \chi_2}{1 + q}, \quad \chi_- = \frac{q\chi_1 - \chi_2}{1 + q}, \quad \delta = \sqrt{1 - 4\eta} \quad \text{with} \quad \eta = \frac{q}{(1 + q)^2}$$

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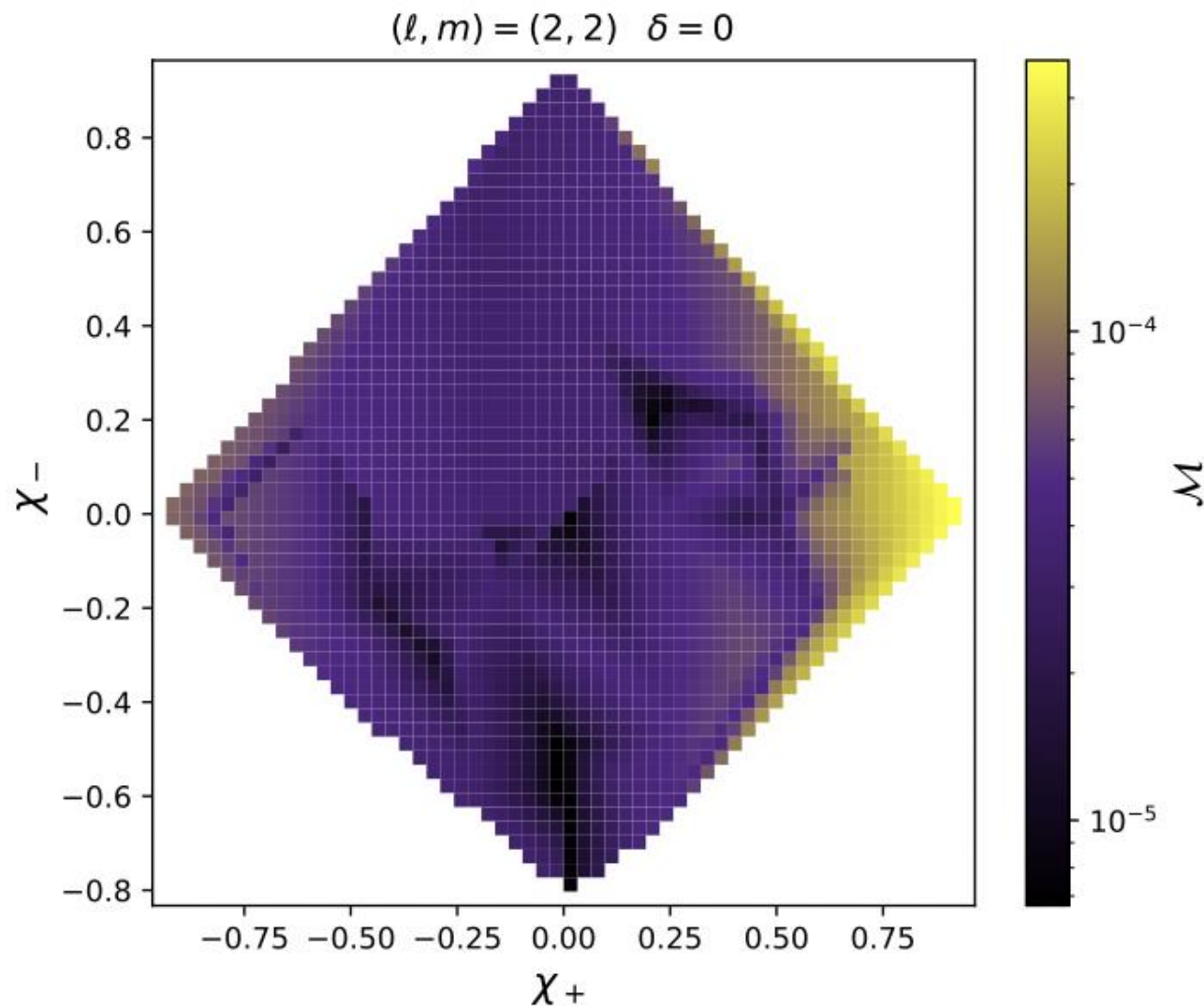
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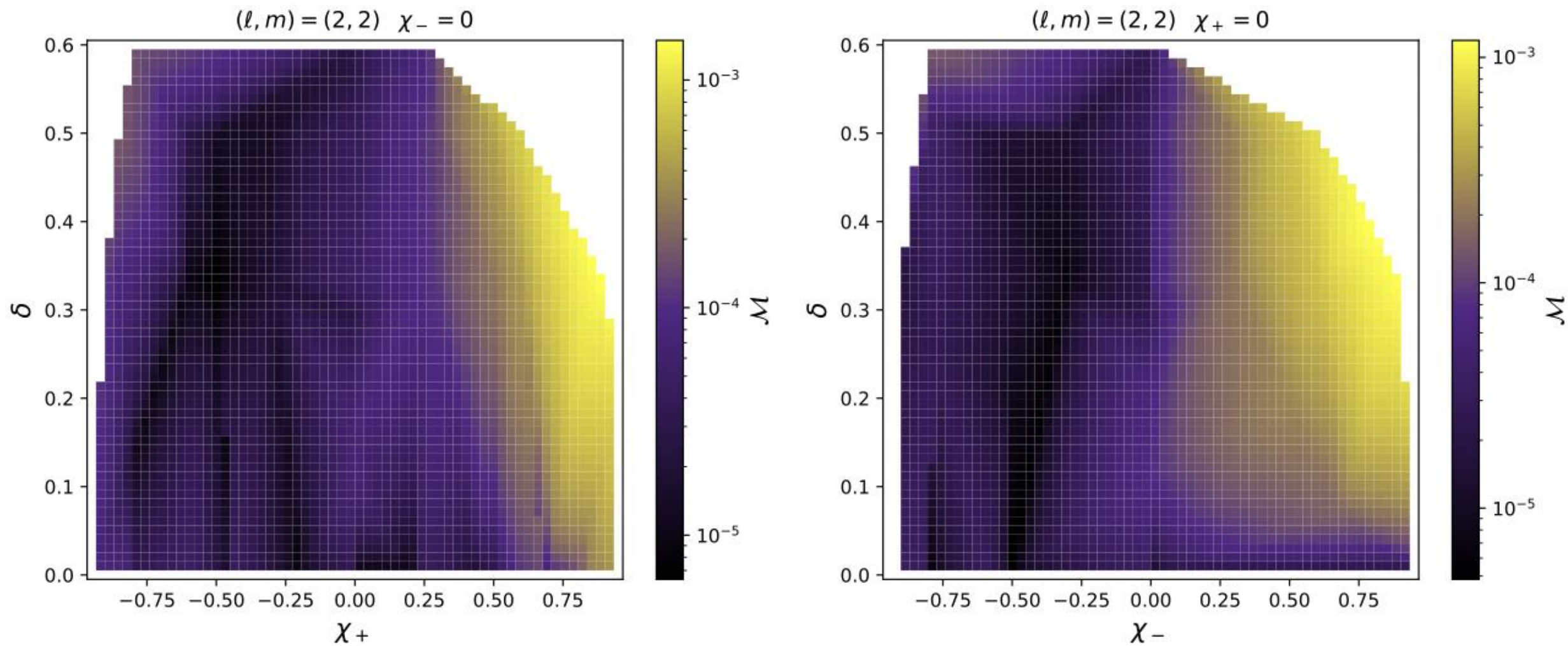
- **Spin aligned**
- **Non eccentric**
- **Comparable mass ($\eta \in [0.15, 0.25]$)**

II. FITTING $A_{\ell m}$ and $p_{\ell m}$ WITH BINARY PARAMETERS

We performed the fit on a subset of ~ 250 simulations and then tested its performance on a completely independent subset of ~ 100 simulations.



II. FITTING $A_{\ell m}$ and $p_{\ell m}$ WITH BINARY PARAMETERS



TAKE-ON MESSAGES

1. Greybody factors can represent a robust gravitational wave observable.
2. Comparable mass mergers can be fitted with GFs.

FUTURE PROSPECTS

1. Understanding the role of spin induced spherical-spheroidal mixing (in preparation)
2. Including precession.
3. Applying the model to real data (in preparation).

Thank you for listening ☺

You can find everything about comparable mass mergers here:
Modeling the frequency-domain ringdown amplitude of comparable-mass mergers with greybody factors, **Romeo Felice Rosato, Sophia Yi, Emanuele Berti and Paolo Pani.**

<https://arxiv.org/pdf/2512.15877>